



TECHNISCHE
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Systematic coarse graining using reversible work potentials

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Computational
Physical Chemistry

Spring school **MULTISCALE SIMULATION OF SOFT MATTER**, Shiraz, Iran, April 9-12, 2018

Outline

- Recap of free energy calculation methods
- Conditional Reversible Work (**CRW**) coarse graining

Monday, April 9

- Applications to soft matter problems
- Dynamically-consistent coarse-grained models

Tuesday, April 10

Multiscale simulations of soft matter

Time

Length

QM model

Atomistic model

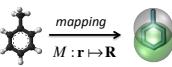
Supra-atomistic model

Linking chemistry and properties on large time and length scales!

Systematic coarse graining

Thermodynamic consistency

$$Q^{CG} = Q^{CG} = \int \exp[-\beta U^{CG}(\mathbf{R})] d\mathbf{R} \quad \beta = (k_B T)^{-1}$$



$$U^{CG}(\mathbf{R}) = -k_B T \ln \left[\int e^{-\beta U^{CG}(\mathbf{r})} \delta(M(\mathbf{r}) - \mathbf{R}) d\mathbf{r} \right]$$

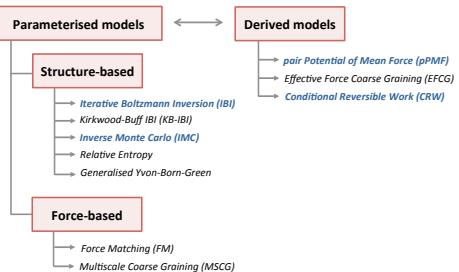
\mathbf{R} : coarse DOFs

\mathbf{r} : atomic DOFs

- Multi-body potential of mean force
- Prohibitively difficult to calculate and represent

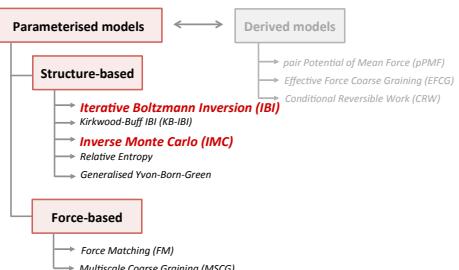
$$\begin{aligned} U^{CG}(\mathbf{R}) &= \sum_{i \neq j} U_2(\mathbf{R}_i, \mathbf{R}_j) + \sum_{i \neq j \neq k} \sum_{k} U_3(\mathbf{R}_i, \mathbf{R}_j, \mathbf{R}_k) + \dots \\ &= \sum_{i \neq j} U_{eff}(\mathbf{R}_i) \end{aligned}$$

Methods for effective pair potentials



E. Brini et al. *Soft Matter* **9**, 2108 (2013)

Methods for effective pair potentials

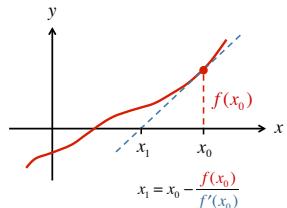


E. Brini et al. *Soft Matter* **9**, 2108 (2013)

Newton's method (Newton-Raphson)

- Finding roots (zeroes) of a function $x : f(x)=0$
 - Initial guess: x_0
 - Iterate:

$$x_{n+1} = x_n - \frac{f(x_n)}{f'(x_n)}$$

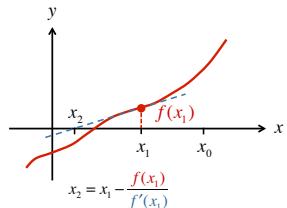


$$x_1 = x_0 - \frac{f(x_0)}{f'(x_0)}$$

Newton's method (Newton-Raphson)

- Finding roots (zeroes) of a function $x : f(x) = 0$
 - Initial guess: x_0
 - Iterate:

$$x_{n+1} = x_n - \frac{f(x_n)}{f'(x_n)}$$

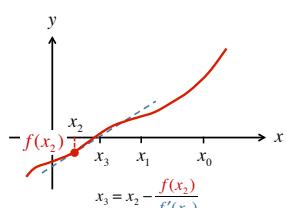


$$x_2 = x_1 - \frac{f(x_1)}{f'(x_1)}$$

Newton's method (Newton-Raphson)

- Finding roots (zeroes) of a function $x : f(x) = 0$
 - Initial guess: x_0

Iterate:



$$x_3 = x_2 - \frac{f(x_2)}{f'(x_2)}$$

Iterative Boltzmann Inversion (IBI)

$$u_{n+1} = u_n + k_B T \ln \left[g_n / g^{ref} \right] \quad n = 0, 1, 2, \dots$$

$F(u_n) = g_n$ Operator F maps u on g

Reith, Putz, Müller-Plathe, JCC (2003)

Start with the potential of mean force: $u_0 = -k_B T \ln g^{ref}$

Setting $Y(u) = -k_B T \ln F(u)$ this is equivalent to

$$u_{n+1} = u_n + (u_0 - Y(u_n)) \quad n = 0, 1, 2, \dots$$

For comparison: The classical Newton method would be

$$u_{n+1} = u_n + Y'(u_n)^{-1} (u_0 - Y(u_n)) \quad n = 0, 1, 2, \dots$$

IBI is a modified Newton method

This shows that IBI is a modified Newton method which uses the approximate derivative

$$Y'(u) = -k_B T \frac{1}{F(u)} F'(u) \approx Id$$

i.e.,

$$F'(u) \approx -\beta F(u) Id$$

Becomes exact when $F(u) = e^{-\beta u}$ (low density limit)

Inverse Monte Carlo is an exact Newton method

Another option is the inverse Monte Carlo iteration

$$u_{n+1} = u_n + K^{-1} (g^{ref} - F(u_n)) \quad n = 0, 1, 2, \dots$$

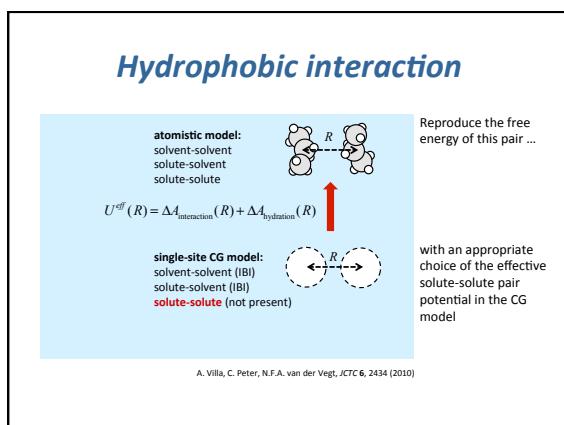
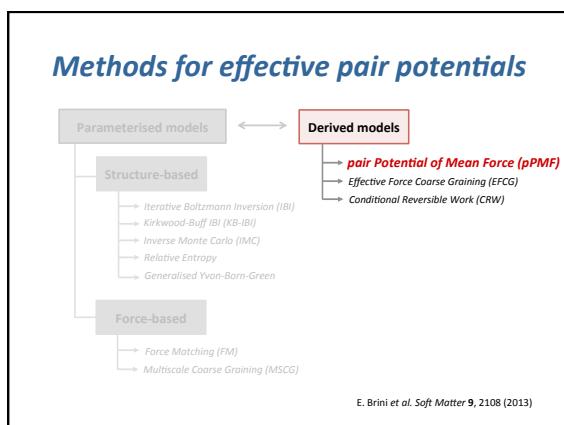
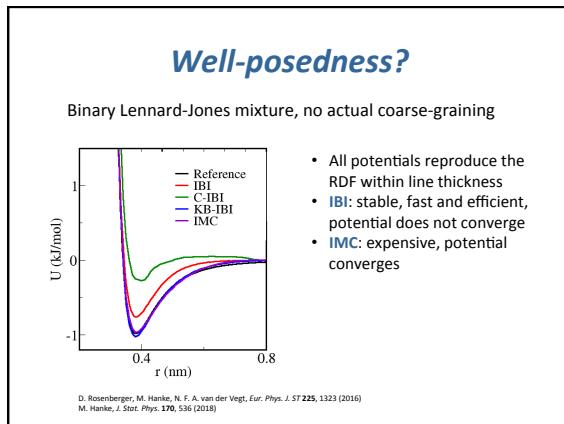
where K is a discretization of the integral operator

$$(Ku)(r) = \int_0^{\infty} k(r, r') u(r') dr'$$

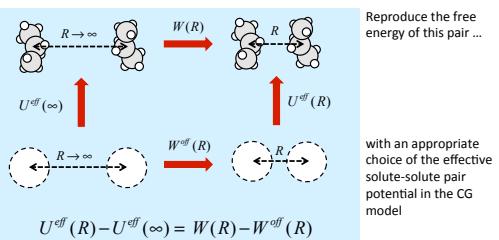
with

$$k(r, r') = -\beta (\langle S(r) S(r') \rangle - \langle S(r) \rangle \langle S(r') \rangle)$$

Lyubartsev, Laaksonen, 1995



Pair potential of mean force



Reproduce the free energy of this pair ...

with an appropriate choice of the effective solute-solute pair potential in the CG model

A. Villa, C. Peter, N.F.A. van der Vegt, JCTC 6, 2434 (2010)

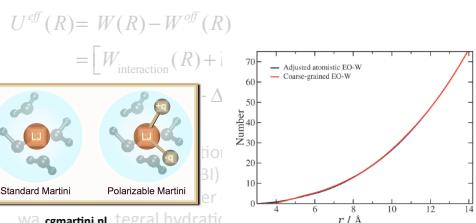
The effective potential

$$\begin{aligned} U^{\text{eff}}(R) &= W(R) - W^{\text{off}}(R) \\ &= [W_{\text{interaction}}(R) + W_{\text{solvent}}(R)] - W_{\text{solvent}}^{\text{CG}}(R) \\ &= W_{\text{interaction}}(R) + \Delta W_{\text{solvent}}(R) \end{aligned}$$

- Can be used in combination with any existing CG solvent model (not necessarily IBI)
- For example: water-water (MARTINI, 4:1 mapping), solute-water (IBI, integral hydration number)
- Tradeoff: structure versus thermodynamics

A. Villa, C. Peter, N.F.A. van der Vegt, JCTC 6, 2434 (2010)

The effective potential



Analogy: single-iteration IBI

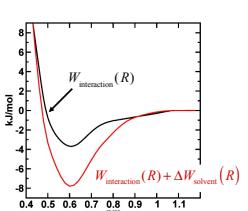
$$u_{n+1} = u_n + (W - Y(u_n)) \quad n = 0, 1, 2, \dots$$

$$W = -k_B T \ln g^{\text{ref}}$$

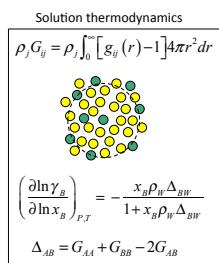
- perform a single IBI iteration on the solute-solute interaction using fixed (pre-determined) solute-water and water-water IBI interactions
- choose $u_0 = 0$

► $u_1 = W - Y(u_0) = W - W^{\text{off}}$
 $= W_{\text{interaction}}(R) + \Delta W_{\text{solvent}}(R)$

Single site benzene/water model



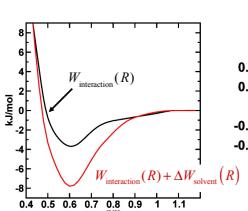
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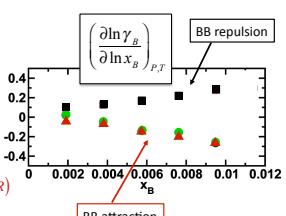
J. Milletti, D. Nayar, N.F.A. van der Vegt, J. Phys. Chem. B (2018)

P. Ganguly, N.F.A. van der Vegt, JCTC 9, 1347 (2013)

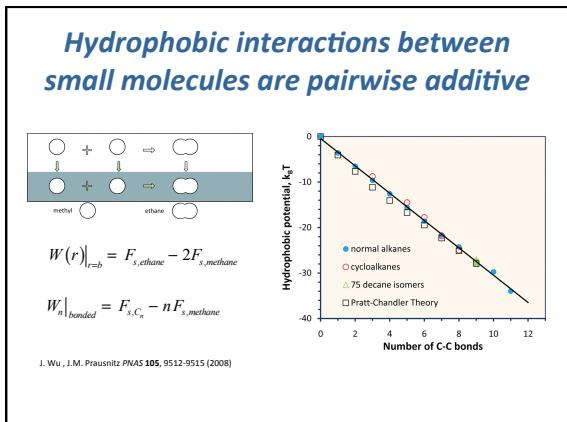
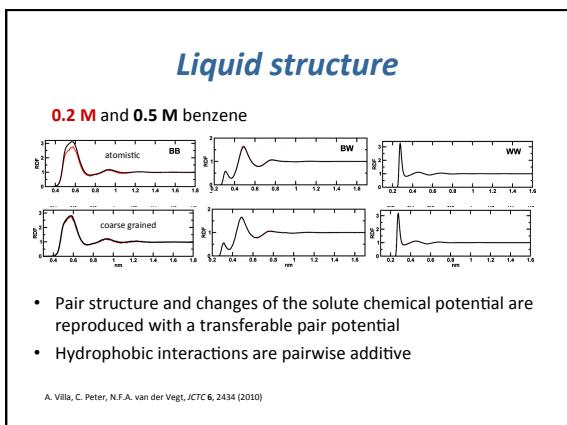
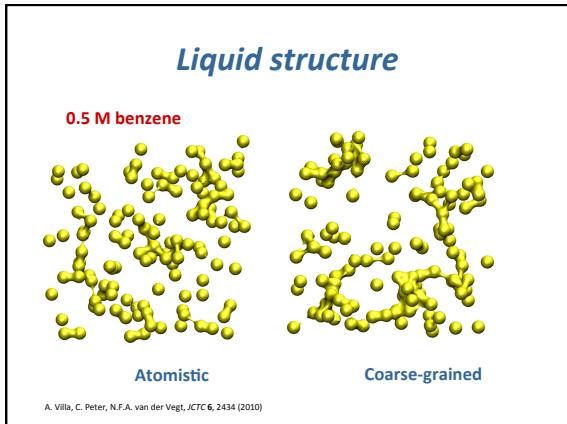
Single site benzene/water model

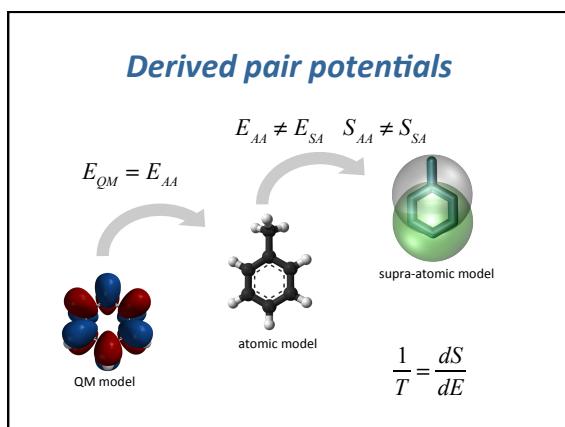
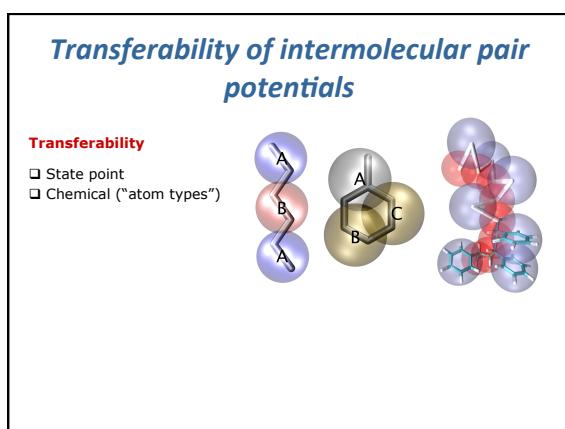
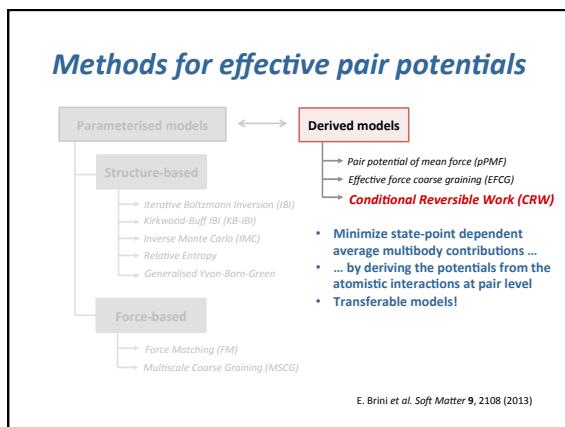


A. Villa, C. Peter, N.F.A. van der Vegt, JCTC 6, 2434 (2010)



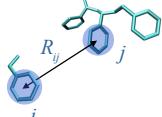
Hydrophobic interactions (in this system)
are pairwise additive





Pair interaction free energy

- No a priori knowledge about the functional dependence on R_{ij}
- Account for entropy of lost degrees of freedom
- Free energy calculations $\Rightarrow \Delta A_g(R_{ij})$
- Use this quantity as a pair potential
- The **Conditional Reversible Work (CRW)** method

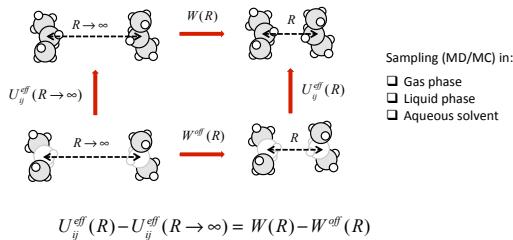


E. Brini et al. PCCP (2011); PCCP (2012); J. Chem. Phys. (2012)
G. Deichmann et al. JCP 141, 224109 (2014)
V. Ardonha et al. JCP 143, 244135 (2015)

Y. Wang, W.G. Noid, P. Liu, G.A. Voth, PCCP (2009)
D. Fritz, V. Harmandaris, K. Kremer, N.F.A. van der Vegt, Macromolecules (2009)
N. Zachariopoulos, N. Vergadou, D.N. Theodorou, J. Chem. Phys. (2005)
J.D. McCoy, J.G. Curro, Macromolecules (1998)

Conditional Reversible Work (CRW)

Constrained sampling (**Cycle** approach)



Sampling (MD/MC) in:
 Gas phase
 Liquid phase
 Aqueous solvent

Conditional Reversible Work (CRW)

Unconstrained sampling (**FEP** approach)

- Simulation of state **1**: liquid phase, full Hamiltonian
- Statistics: Averaging over all pairs

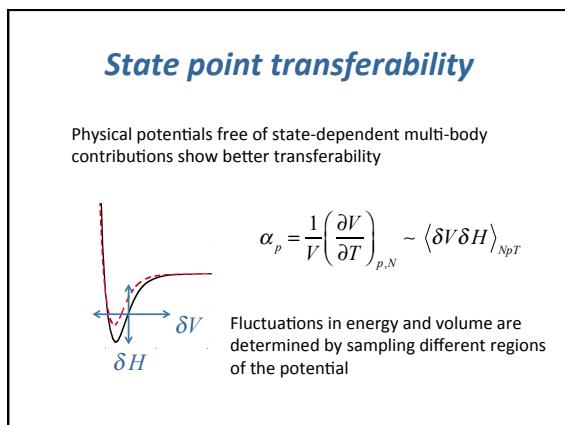
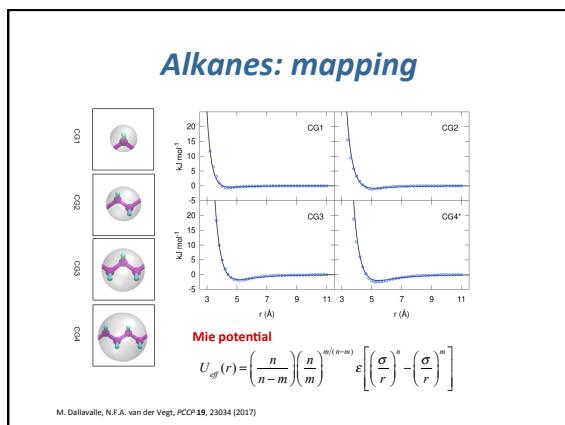
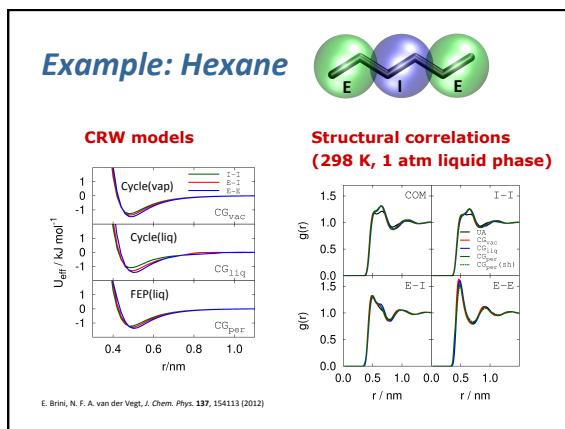
$$\begin{aligned} U_{ij}^{eff}(R) &= \Delta A_{\text{1} \rightarrow \text{1}}(R) \\ &= k_B T \ln \left\langle e^{\beta E_j(R)} \right\rangle_1 \\ &= k_B T \ln \left[\int_{-\infty}^{\infty} P_j(E; R) \exp(-\beta E) dE \right] \end{aligned}$$

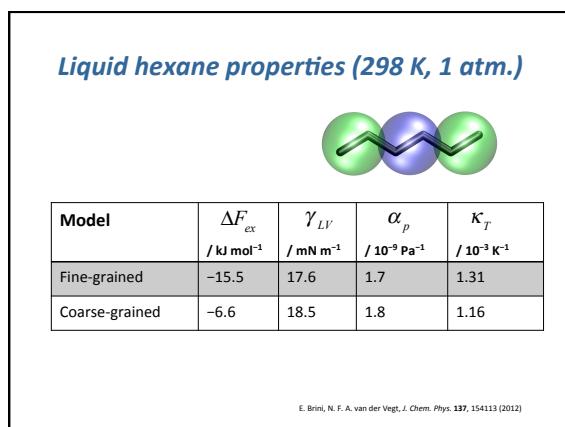
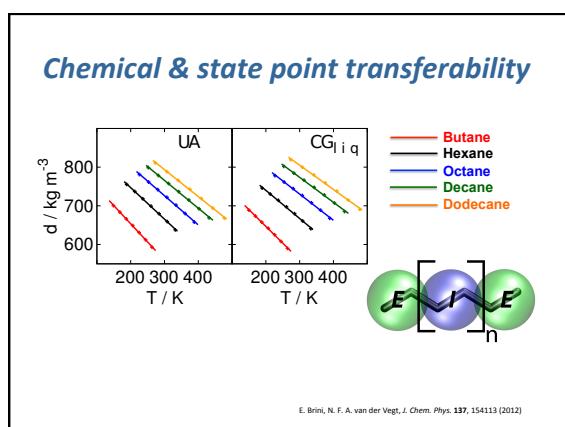
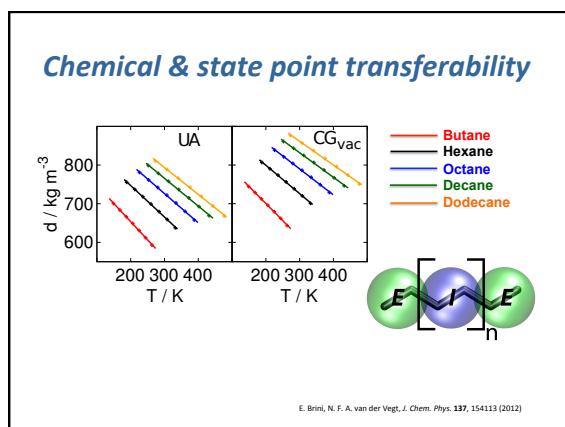
Reference ensemble (state **1**)

$\Delta A_{\text{1} \rightarrow \text{1}}(R) = U_{ij}^{eff}(R)$

Perturbed ensemble (state **0**)

E. Brini, N. F. A. van der Vegt, J. Chem. Phys. 137, 154113 (2012)





Transferability: Material surfaces

D. Fritz et al. *Macromolecules* **42**, 7579 (2009)

V. Marcon, N.F.A. van der Vegt, *Soft Matter* **8**, 5585 (2012)

Cooling

